



Hidden On-Shell Mediators for the Galactic Center Excess

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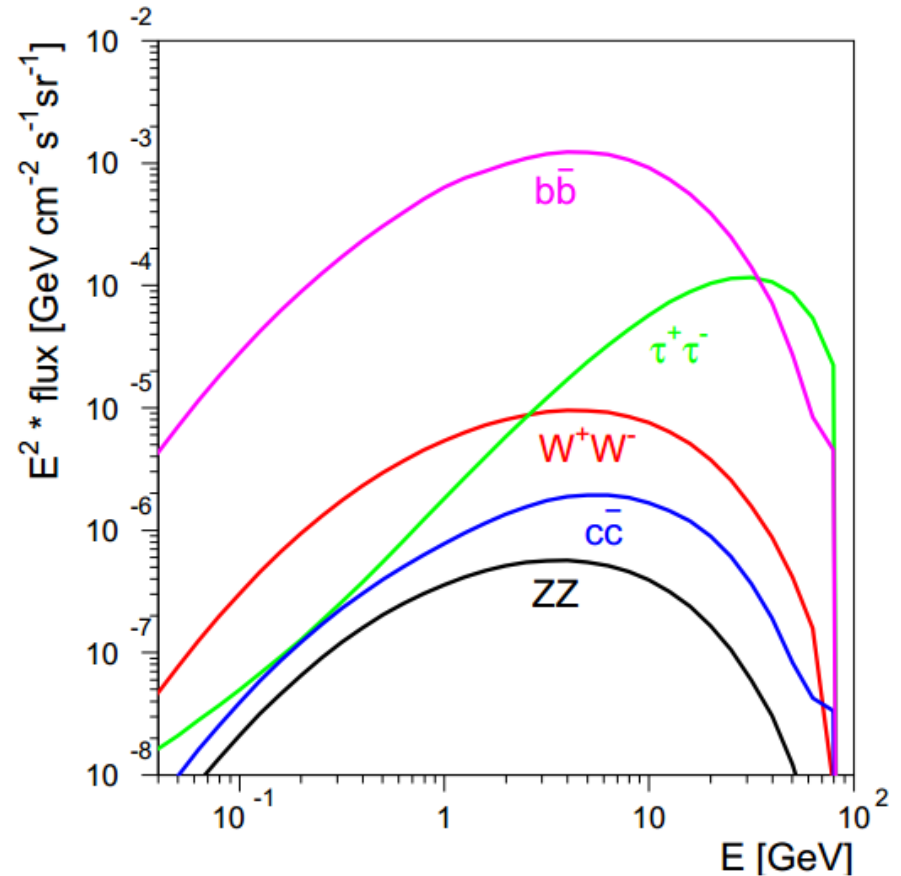
Work done with:
M. Abdullah, A. DiFranzo, A. Rajaraman,
T. M.P. Tait, P. Tanedo
[arXiv:1404.6528 \[hep-ph\]](https://arxiv.org/abs/1404.6528)

Outline

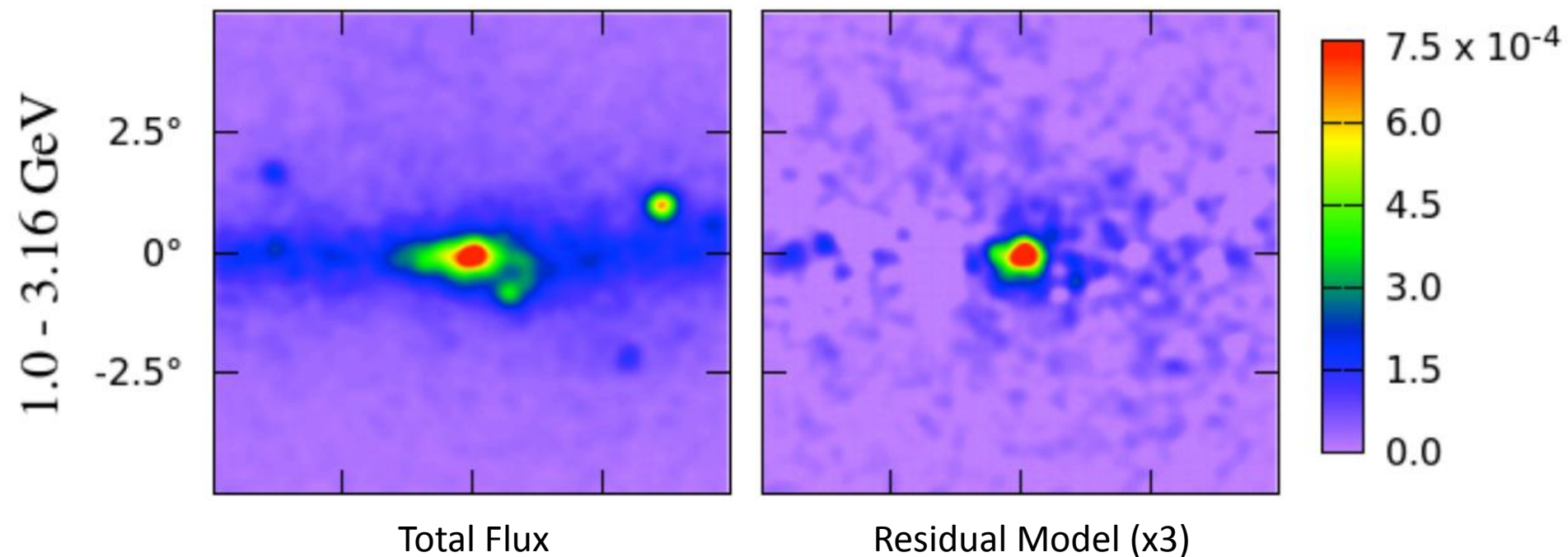
- What is the galactic center excess?
- Particle physics models for the galactic center
 - Mechanism to model the excess via on-shell mediators
- Parameter space for the galactic center
- Bounds on hidden, on-shell mediators
- Opportunities for future work

Indirect Detection

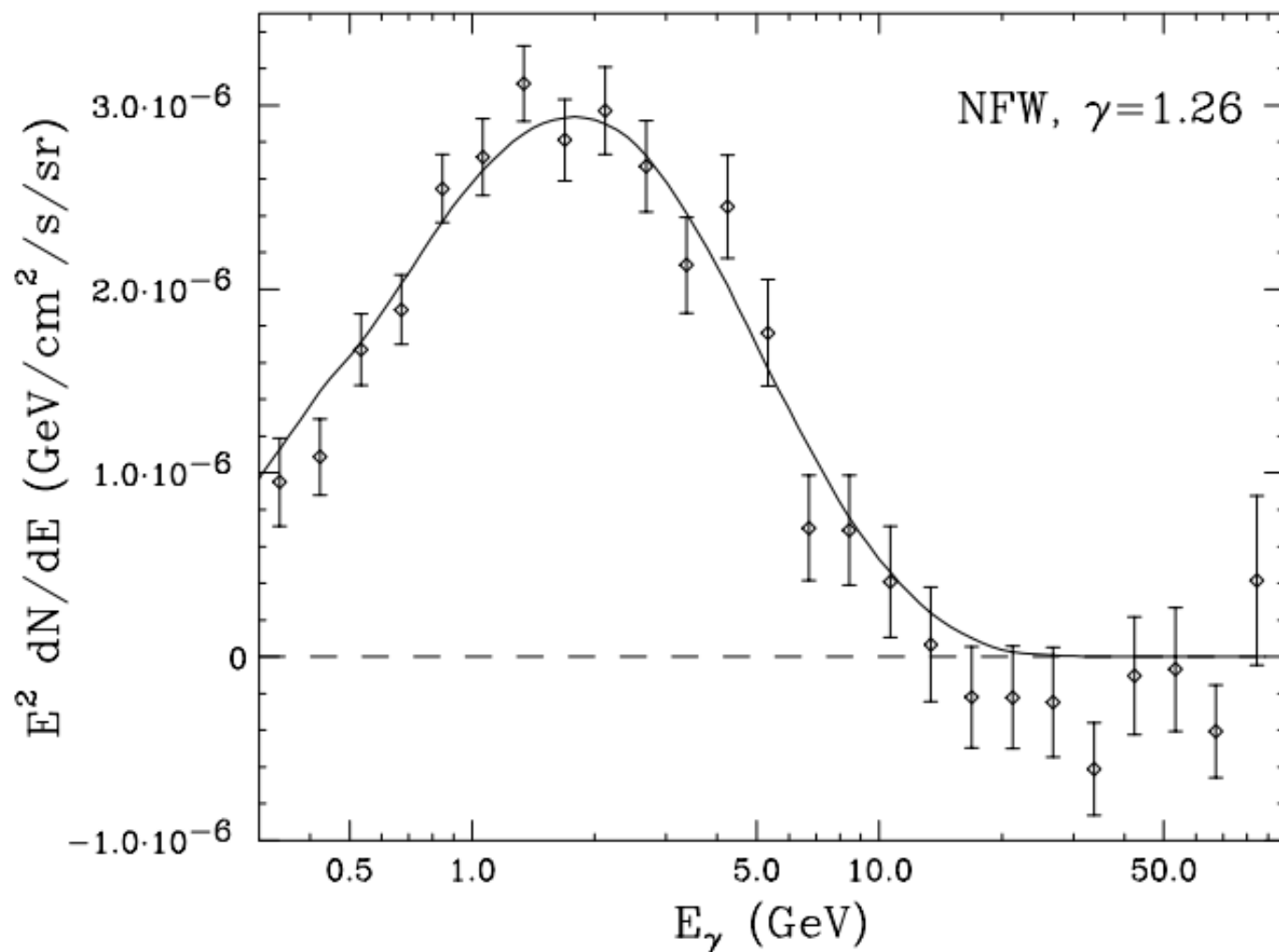
- Dark Matter may annihilate or decay into standard model particles
- Eventually, these particles become protons, electrons, neutrinos, and photons
- The photons can reach the Earth and are reliably detected.



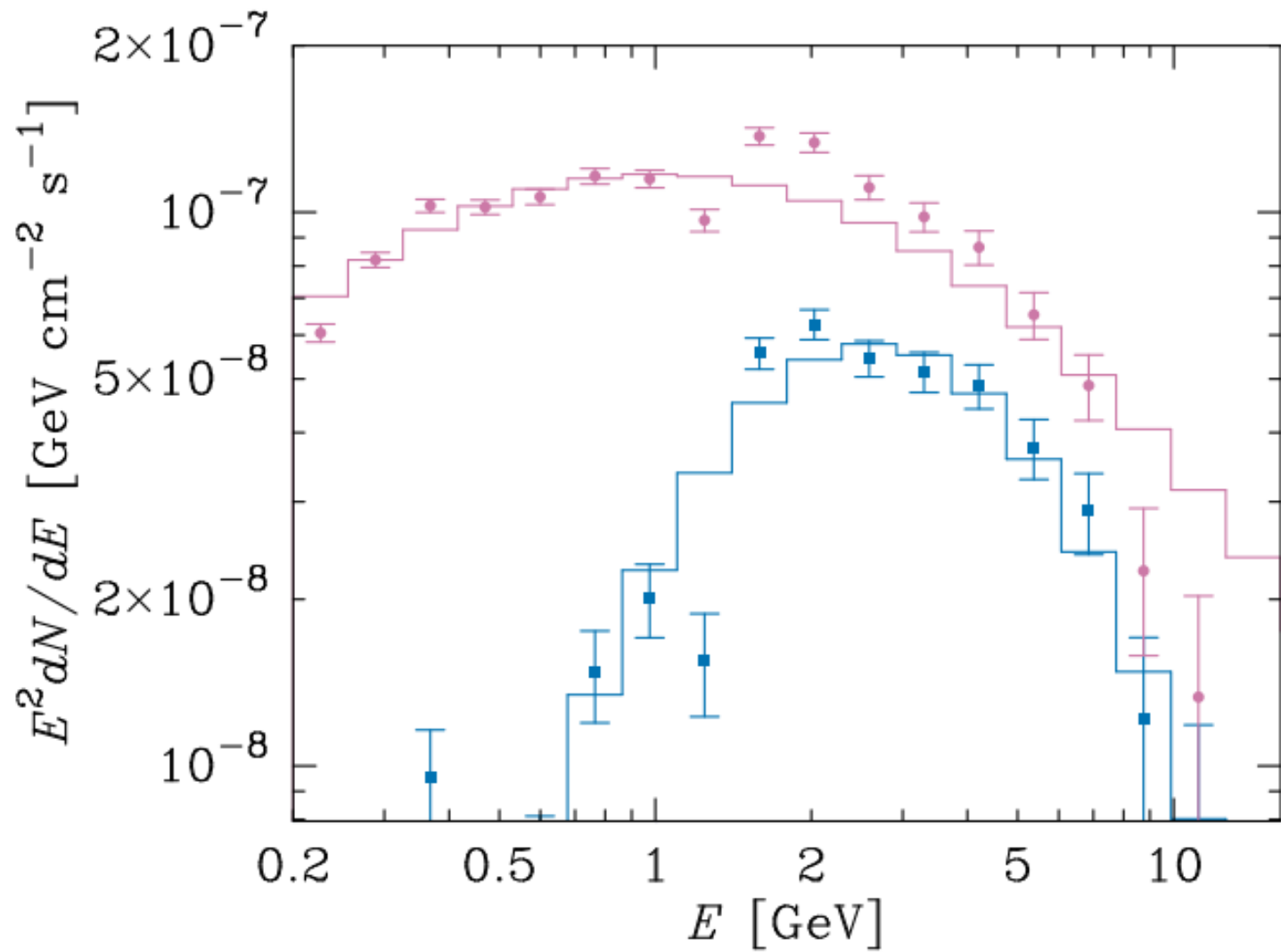
What is the Galactic Center Excess?



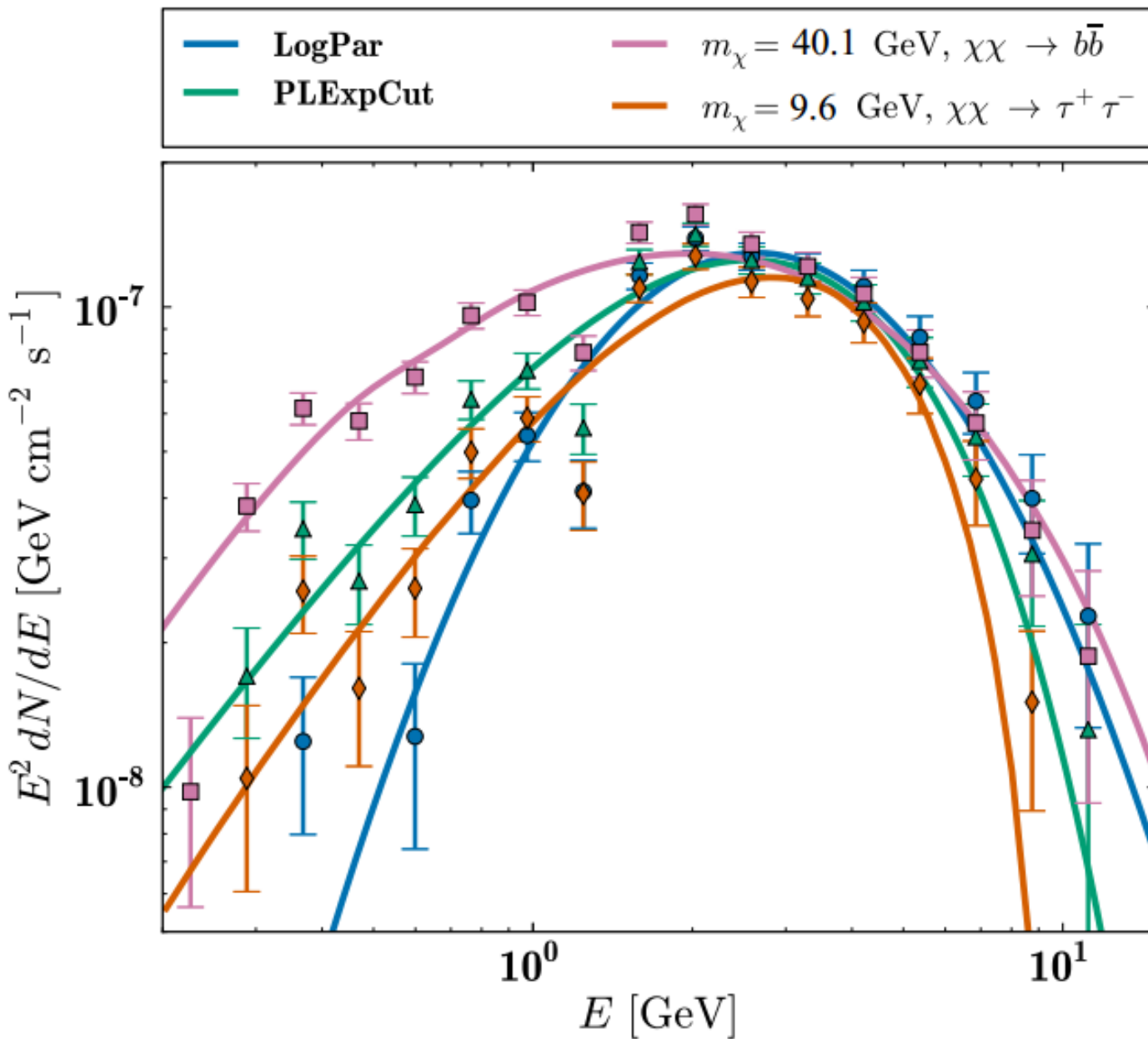
$$\rho(r) = \rho_0 \frac{(r/r_s)^{-\gamma}}{(1 + r/r_s)^{3-\gamma}}$$



$$\sigma v = 1.7 \times 10^{-26} \text{ cm}^3/\text{s}$$



$$\sigma v = 5.1 \times 10^{-26} \text{ cm}^3/\text{s}$$

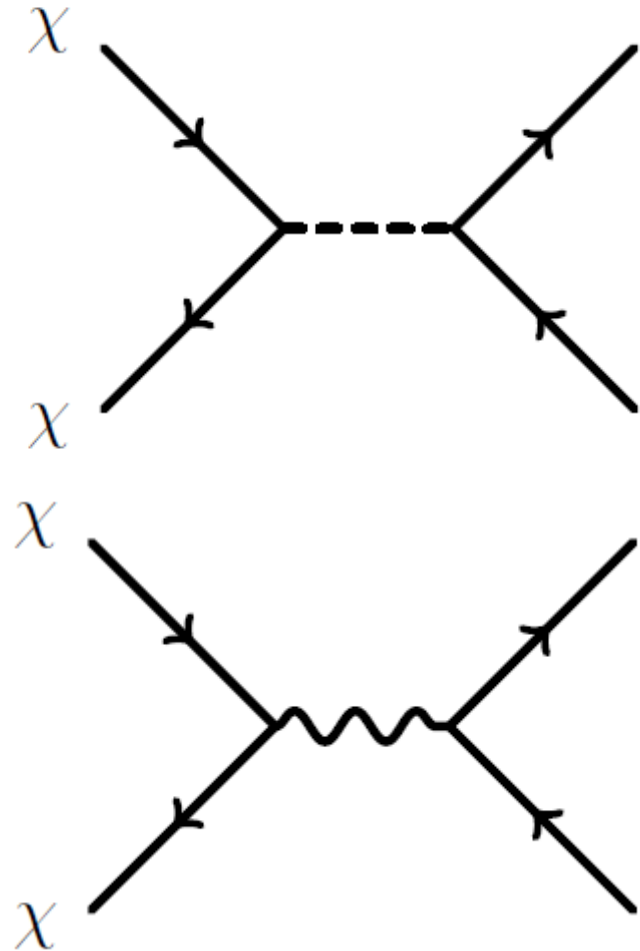


Particle Physics Models for the Galactic Center

- Inputs from the galactic center:
 - Value of σv is suggestive of thermal WIMP
 - S-wave DM annihilation
 - Annihilation should produce τ or b particles
 - The resulting τ or b should have a particular energy to reproduce the observed spectrum

Two Body Annihilation

- Two WIMPs meet and become a pair of SM particles
- Both resulting particles have the energy of the WIMP mass in the center of mass frame
- Since WIMPs are non relativistic, that frame is typically the galactic frame



Two Body Annihilation has been extensively explored

- Effective Interactions
 - A. Alves *et al.* arXiv:1403.5027 [hep-ph]
- Simplified Models
 - A. Berlin *et al.* arXiv:1404.0022 [hep-ph]
 - C. Boehm *et al.* arXiv:1401.6458 [hep-ph]
 - E. Izaguirre *et al.* arXiv:1404.2018 [hep-ph]
 - P. Agrawal *et al.* arXiv:1404.1373 [hep-ph]
- UV Models
 - S. Ipek *et al.* arXiv:1404.3716 [hep-ph]

Effective Operators

$$(D1) \quad \frac{m_q}{\Lambda^3} \bar{\chi} \chi \bar{q} q$$

$$(D2) \quad \frac{m_q}{\Lambda^3} \bar{\chi} \gamma^5 \chi \bar{q} q$$

$$(D3) \quad \frac{m_q}{\Lambda^3} \bar{\chi} \chi \bar{q} \gamma^5 q$$

$$(D4) \quad \frac{m_q}{\Lambda^3} \bar{\chi} \gamma^5 \chi \bar{q} \gamma^5 q$$

$$(D5) \quad \frac{1}{\Lambda^2} \bar{\chi} \gamma^\mu \chi \bar{q} \gamma_\mu q$$

$$(D6) \quad \frac{1}{\Lambda^2} \bar{\chi} \gamma^\mu \gamma^5 \chi \bar{q} \gamma_\mu q$$

$$(D7) \quad \frac{1}{\Lambda^2} \bar{\chi} \gamma^\mu \chi \bar{q} \gamma_\mu \gamma^5 q$$

$$(D8) \quad \frac{1}{\Lambda^2} \bar{\chi} \gamma^\mu \gamma^5 \chi \bar{q} \gamma_\mu \gamma^5 q$$

$$(D9) \quad \frac{1}{\Lambda^2} \bar{\chi} \sigma^{\mu\nu} \chi \bar{q} \sigma_{\mu\nu} q$$

$$(D10) \quad \frac{1}{\Lambda^2} \epsilon^{\mu\nu\alpha\beta} \bar{\chi} \sigma_{\mu\nu} \chi \bar{q} \sigma_{\alpha\beta} q$$

$$(D11) \quad \frac{\alpha_s}{4\Lambda^3} \bar{\chi} \chi (G_{\mu\nu}^a)^2$$

$$(D12) \quad \frac{\alpha_s}{4\Lambda^3} \bar{\chi} \gamma^5 \chi (G_{\mu\nu}^a)^2$$

$$(D13) \quad \frac{\alpha_s}{4\Lambda^3} \bar{\chi} \chi G_{\mu\nu}^a \tilde{G}^{a,\mu\nu}$$

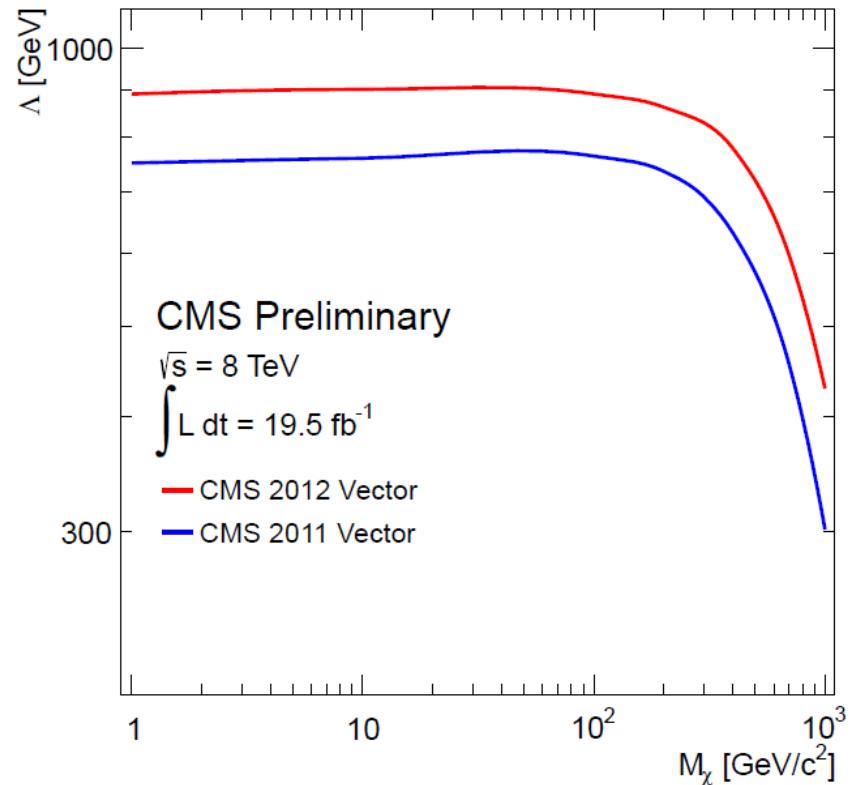
$$(D14) \quad \frac{\alpha_s}{4\Lambda^3} \bar{\chi} \gamma^5 \chi G_{\mu\nu}^a \tilde{G}^{a,\mu\nu}$$

Effective Operators

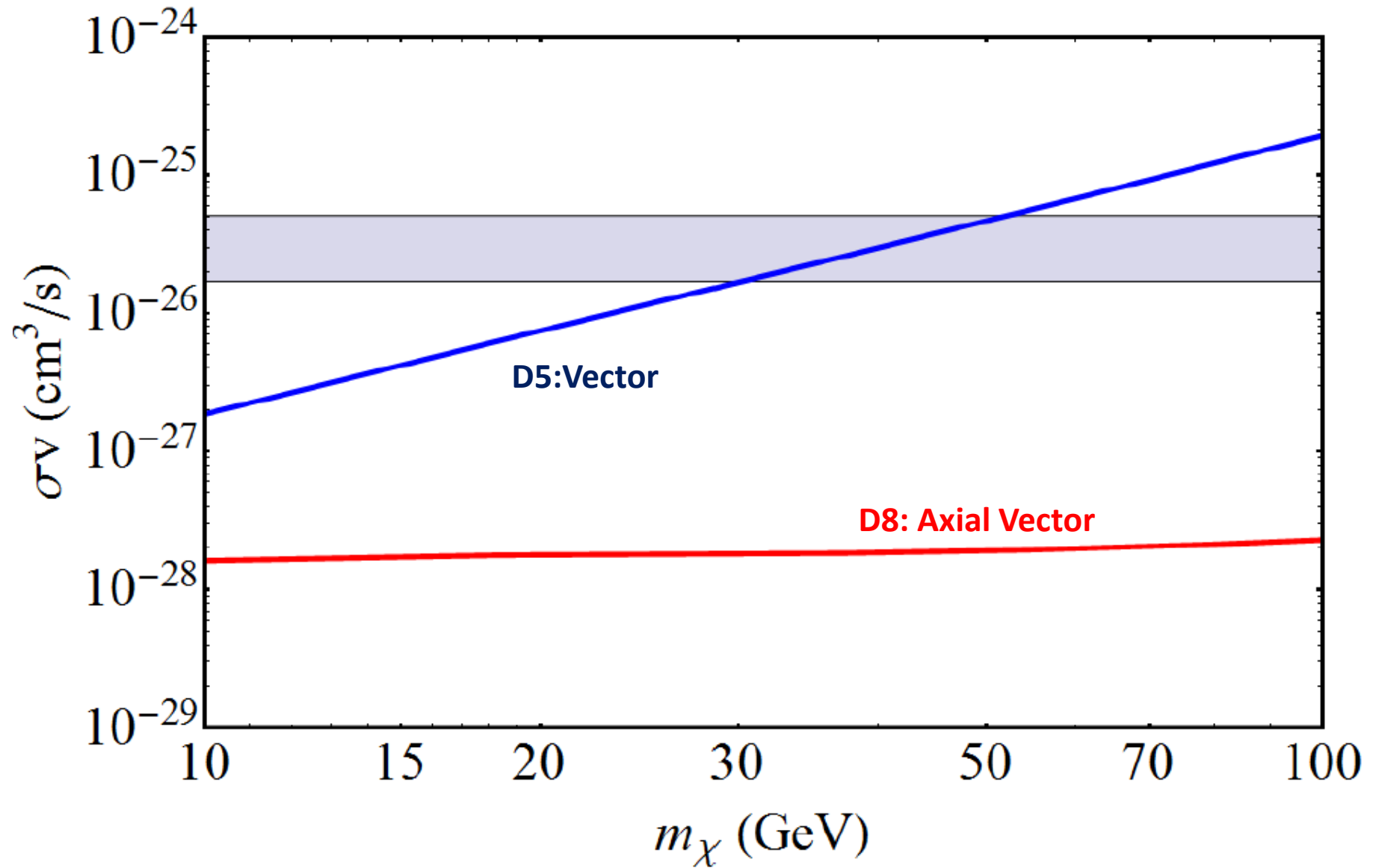
(D1)	$\frac{m_q}{\Lambda^3} \bar{\chi} \chi \bar{q} q$	(D8)	$\frac{1}{\Lambda^2} \bar{\chi} \gamma^\mu \gamma^5 \chi \bar{q} \gamma_\mu \gamma^5 q$
(D2)	$\frac{m_q}{\Lambda^3} \bar{\chi} \gamma^5 \chi \bar{q} q$	(D9)	$\frac{1}{\Lambda^2} \bar{\chi} \sigma^{\mu\nu} \chi \bar{q} \sigma_{\mu\nu} q$
(D3)	$\frac{m_q}{\Lambda^3} \bar{\chi} \chi \bar{q} \gamma^5 q$	(D10)	$\frac{1}{\Lambda^2} \epsilon^{\mu\nu\alpha\beta} \bar{\chi} \sigma_{\mu\nu} \chi \bar{q} \sigma_{\alpha\beta} q$
(D4)	$\frac{m_q}{\Lambda^3} \bar{\chi} \gamma^5 \chi \bar{q} \gamma^5 q$	(D11)	$\frac{\alpha_s}{4\Lambda^3} \bar{\chi} \chi (G_{\mu\nu}^a)^2$
(D5)	$\frac{1}{\Lambda^2} \bar{\chi} \gamma^\mu \chi \bar{q} \gamma_\mu q$	(D12)	$\frac{\alpha_s}{4\Lambda^3} \bar{\chi} \gamma^5 \chi (G_{\mu\nu}^a)^2$
(D6)	$\frac{1}{\Lambda^2} \bar{\chi} \gamma^\mu \gamma^5 \chi \bar{q} \gamma_\mu q$	(D13)	$\frac{\alpha_s}{4\Lambda^3} \bar{\chi} \chi G_{\mu\nu}^a \tilde{G}^{a,\mu\nu}$
(D7)	$\frac{1}{\Lambda^2} \bar{\chi} \gamma^\mu \chi \bar{q} \gamma_\mu \gamma^5 q$	(D14)	$\frac{\alpha_s}{4\Lambda^3} \bar{\chi} \gamma^5 \chi G_{\mu\nu}^a \tilde{G}^{a,\mu\nu}$

CMS Monojet Search

- Attempt to produce a WIMP pair via quark annihilation
- Detectable events occur when the pair is recoiling off QCD ISR, resulting in a single jet and MET

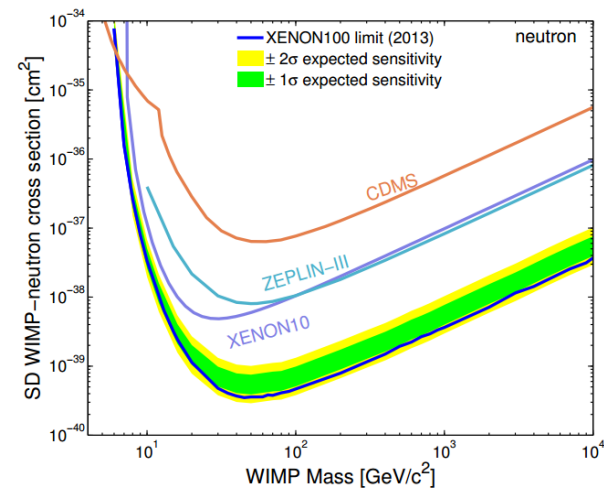
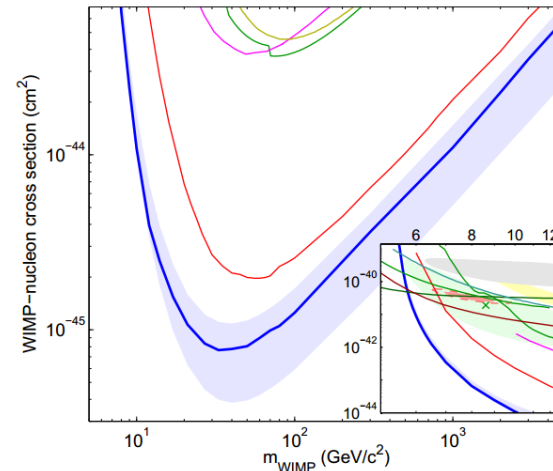


Bounds on Heavy Physics from Colliders



Direct Detection Nucleon Scattering

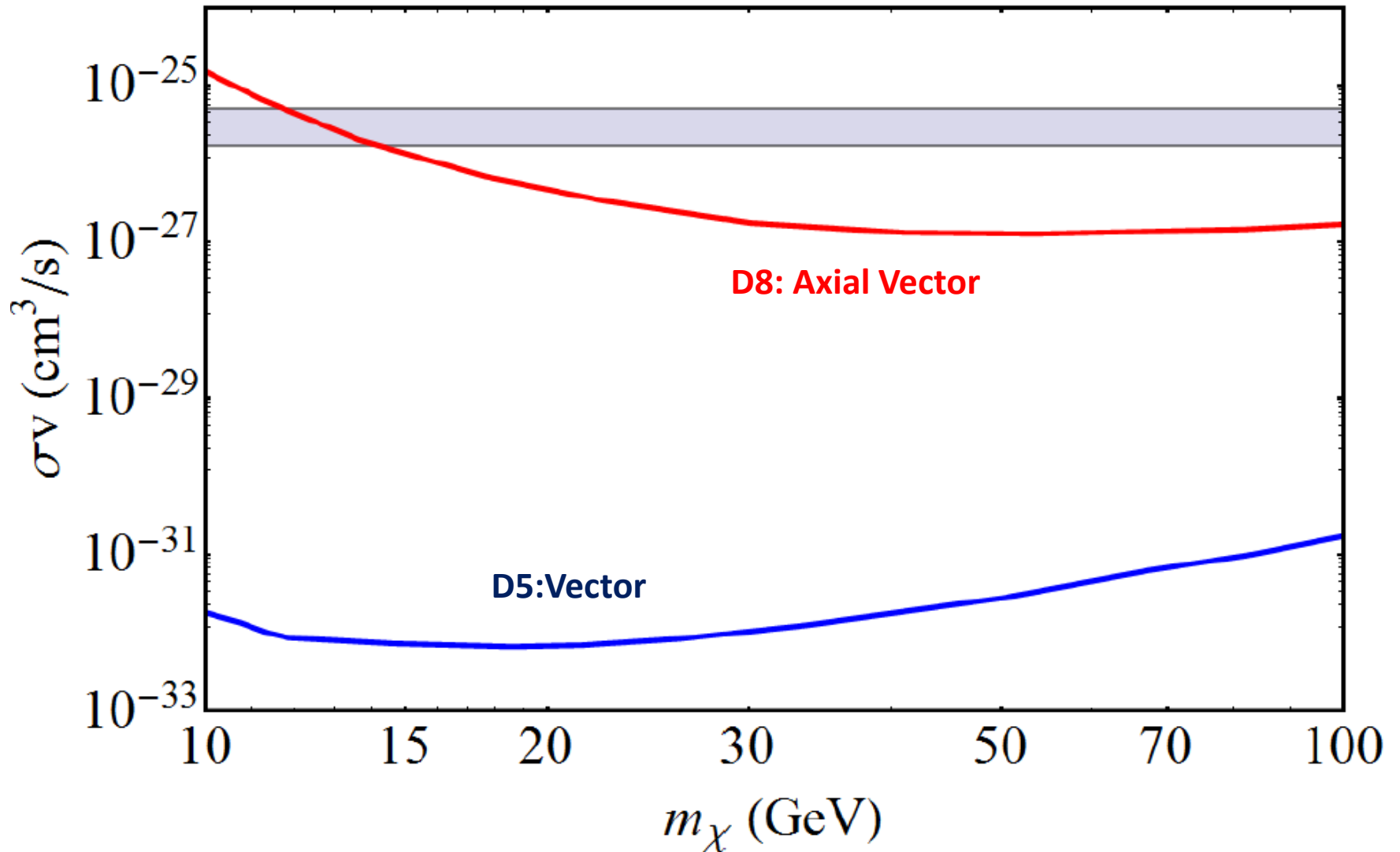
- Attempt to observe WIMP nucleon scattering
- This scattering may or may not be independent of the nucleus spin



D. S. Akerib *et al.* LUX Collaboration, "First results from the LUX dark matter experiment at the Sanford Underground Research Facility," arXiv:1310.8214 [astro-ph.CO].

E. Aprile *et al.* XENON100 Collaboration, "Limits on spin-dependent WIMP-nucleon cross sections from 225 live days of XENON100 data," Phys. Rev. Lett. **111**, no. 2, 021301 (2013) [arXiv:1301.6620 [astro-ph.CO]]

Bounds from Direct Detection



Solution: On-shell mediators

$$\mathcal{L}_{int} = \lambda_{DM} V^\mu \bar{\chi} \gamma_\mu (\gamma^5) \chi + \lambda_{SM} V^\mu \bar{q} \gamma_\mu (\gamma^5) q$$

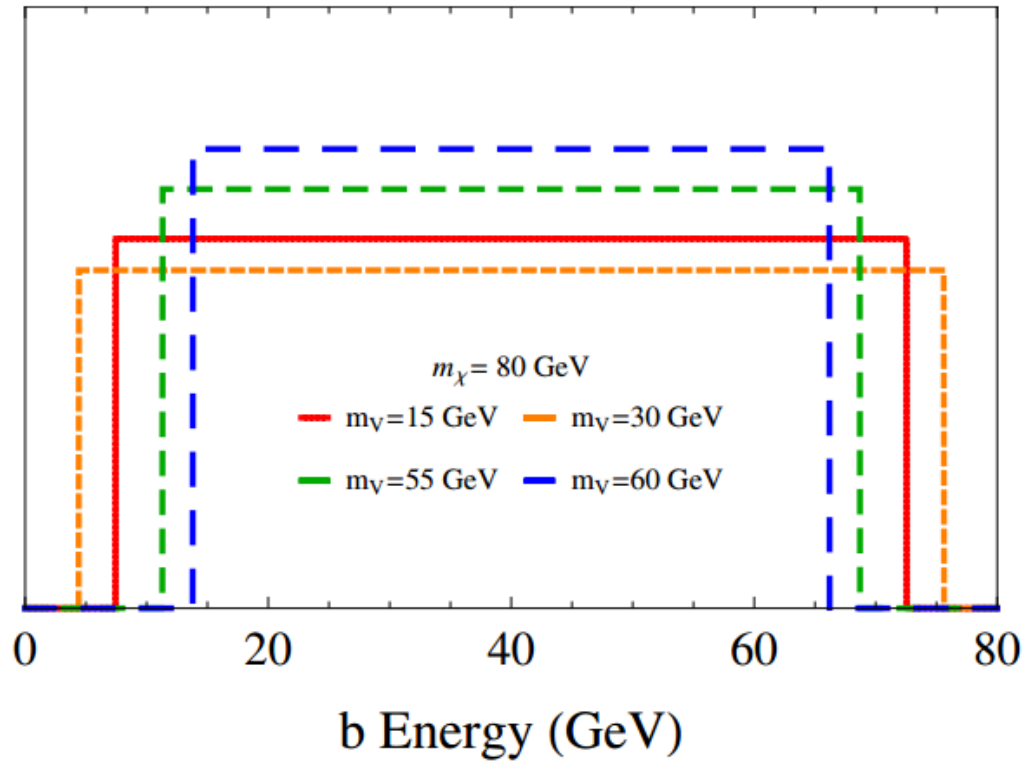
$$\left(\begin{array}{c} \chi \\ \chi \end{array} \right) \left(\begin{array}{c} \text{on shell} \end{array} \right) \sim \lambda_{DM}^2 \gg \sim \lambda_{DM} \lambda_{SM}$$

$$\mathcal{L}_{int} = \lambda_{DM} \phi \bar{\chi} \gamma^5 \chi + \lambda_{SM} \phi \bar{q}_L (\gamma^5) q_R + H.C.$$

$$\left(\begin{array}{c} \chi \\ \chi \end{array} \right) \left(\begin{array}{c} \text{on shell} \end{array} \right) \sim \frac{\lambda_{DM}^3}{\sqrt{4\pi}} \gg \sim \lambda_{DM} \lambda_{SM}$$

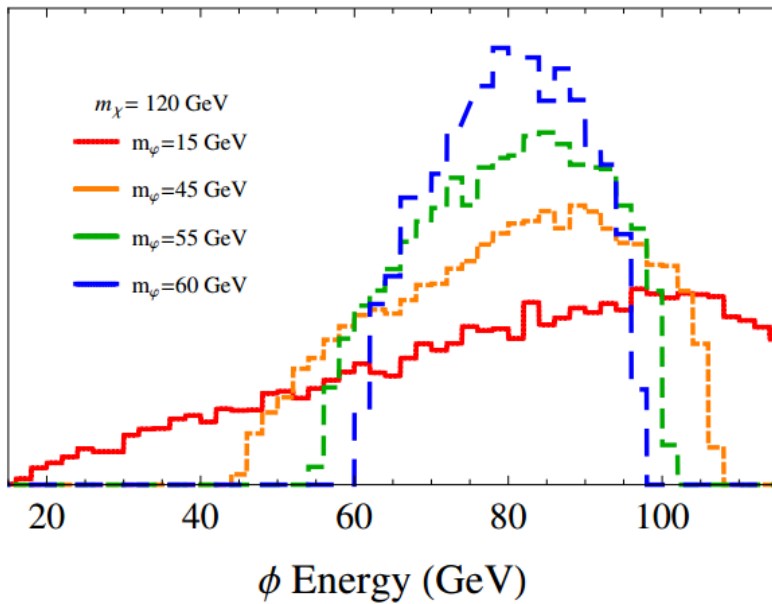
- Inputs from the galactic center:
 - Value of σv is suggestive of thermal WIMP
 - S-wave DM annihilation
 - Annihilation should produce τ or b particles
 - The resulting τ or b should have a particular energy
- Do on-shell mediators satisfy:
 - Allows for WIMPs, thermal relic later
 - 2 vectors or 3 pseudoscalars
 - Resulting particles controlled by mediator couplings
 - Can get 40 GeV b quarks

Spin 1 Mediator

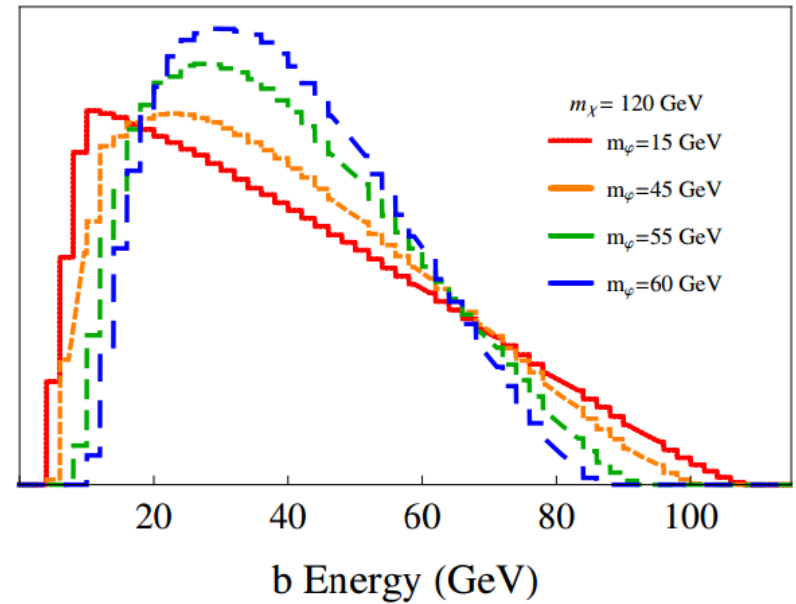


$$\chi\bar{\chi} \rightarrow VV \rightarrow 4b$$

Spin 0 Mediator



$$\chi\bar{\chi} \rightarrow 3\phi$$



$$\chi\bar{\chi} \rightarrow 3\phi \rightarrow 6b$$

Parameter space for the galactic center

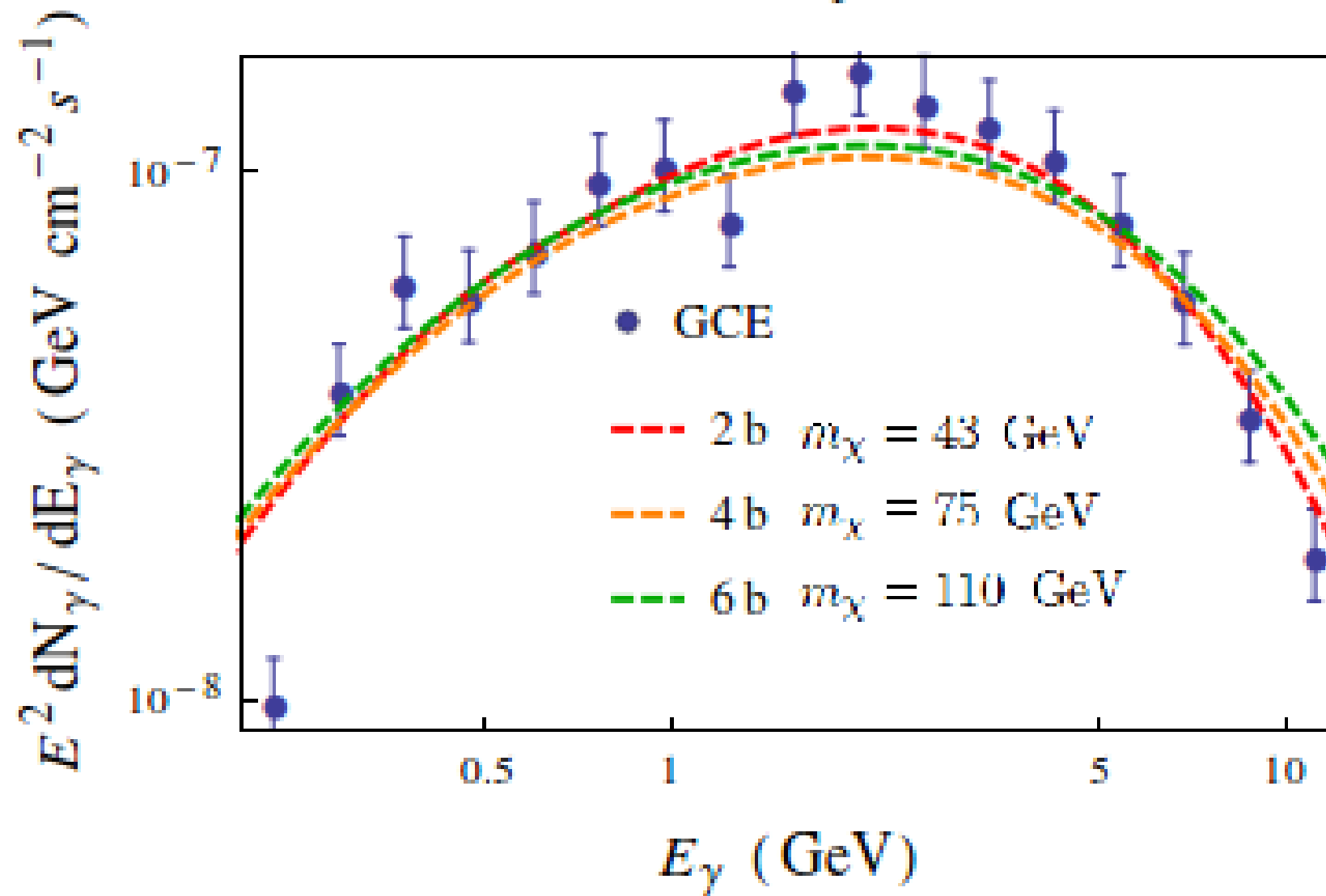
The relevant parameters are:

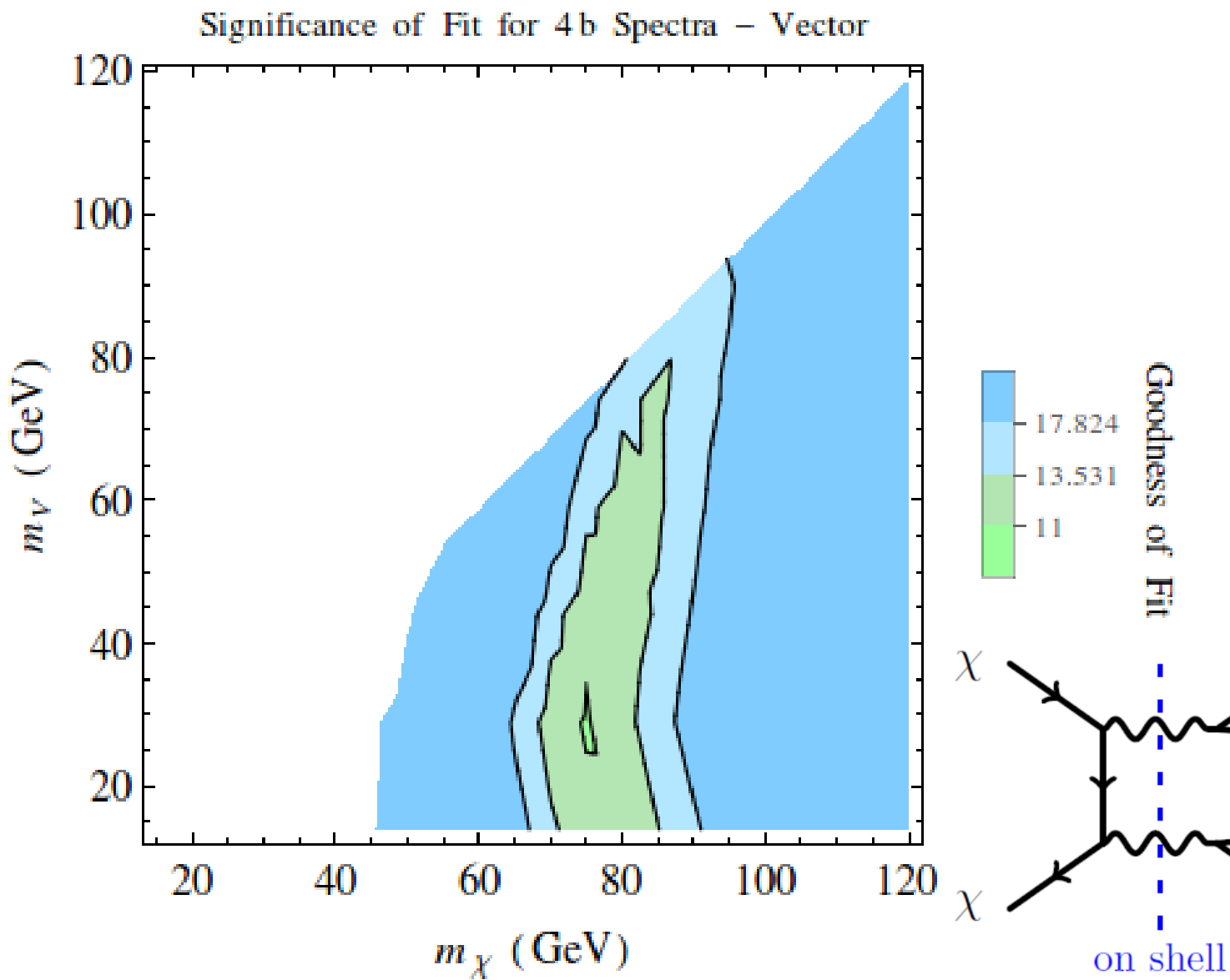
- WIMP mass: controls shape of the spectrum
- Mediator mass: controls shape of the spectrum
- DM Coupling: affects normalization

Since this fit is heavily dependent on the shape, we impose a 20% error on all residuals.

$$\text{goodness of fit} = \sum_i \left(\frac{\log D_i - \log (\lambda_{\text{DM}}^{2n} S_i)}{\log(0.2D_i)} \right)^2$$

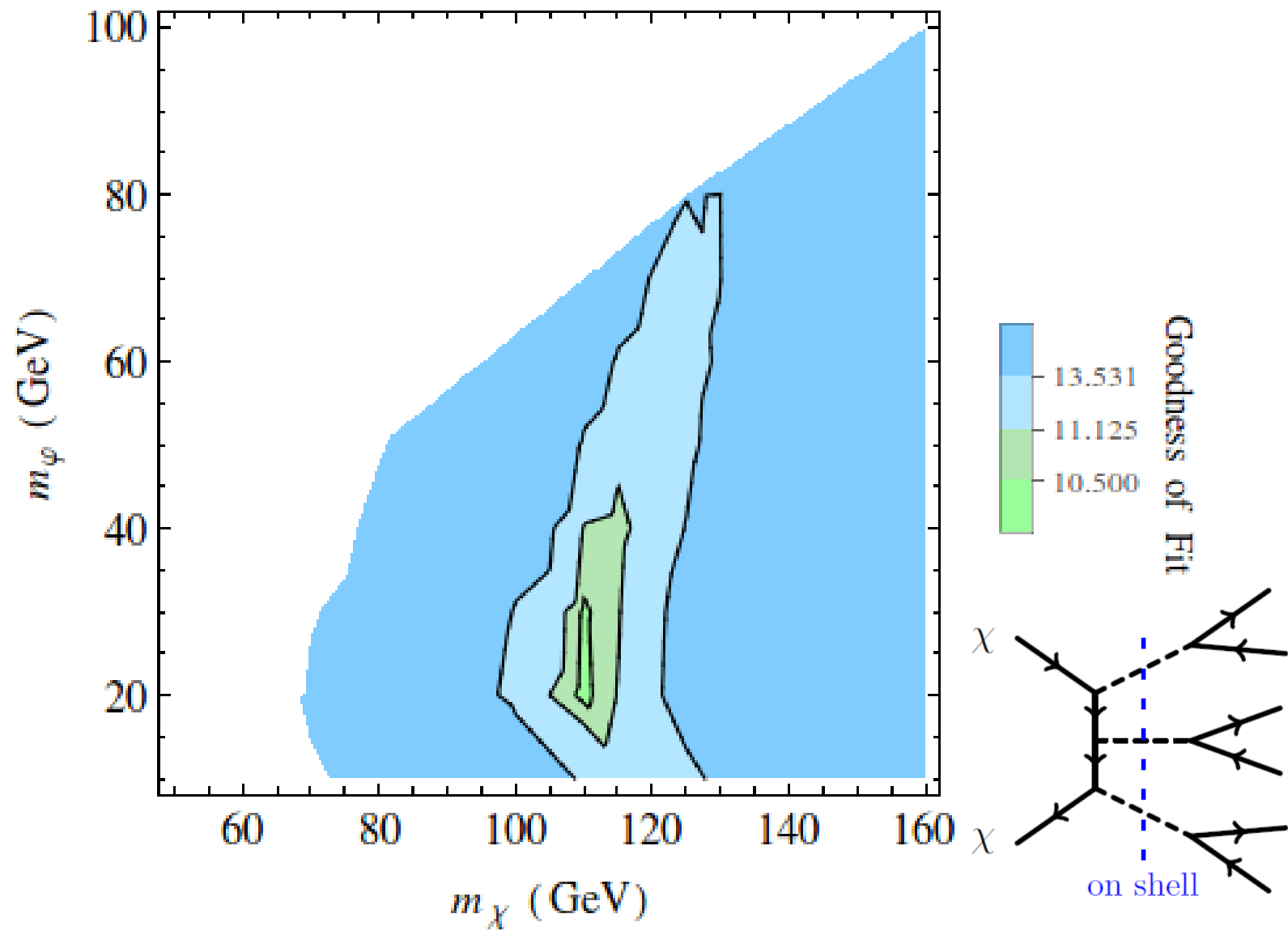
Best Fit Spectra



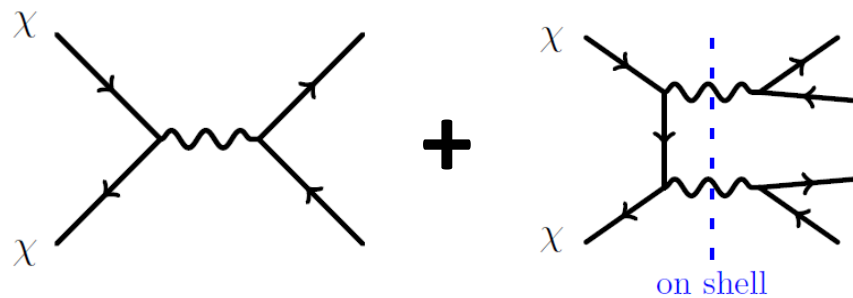


$$\lambda_{\text{DM}} \sim 0.27 - 0.44$$

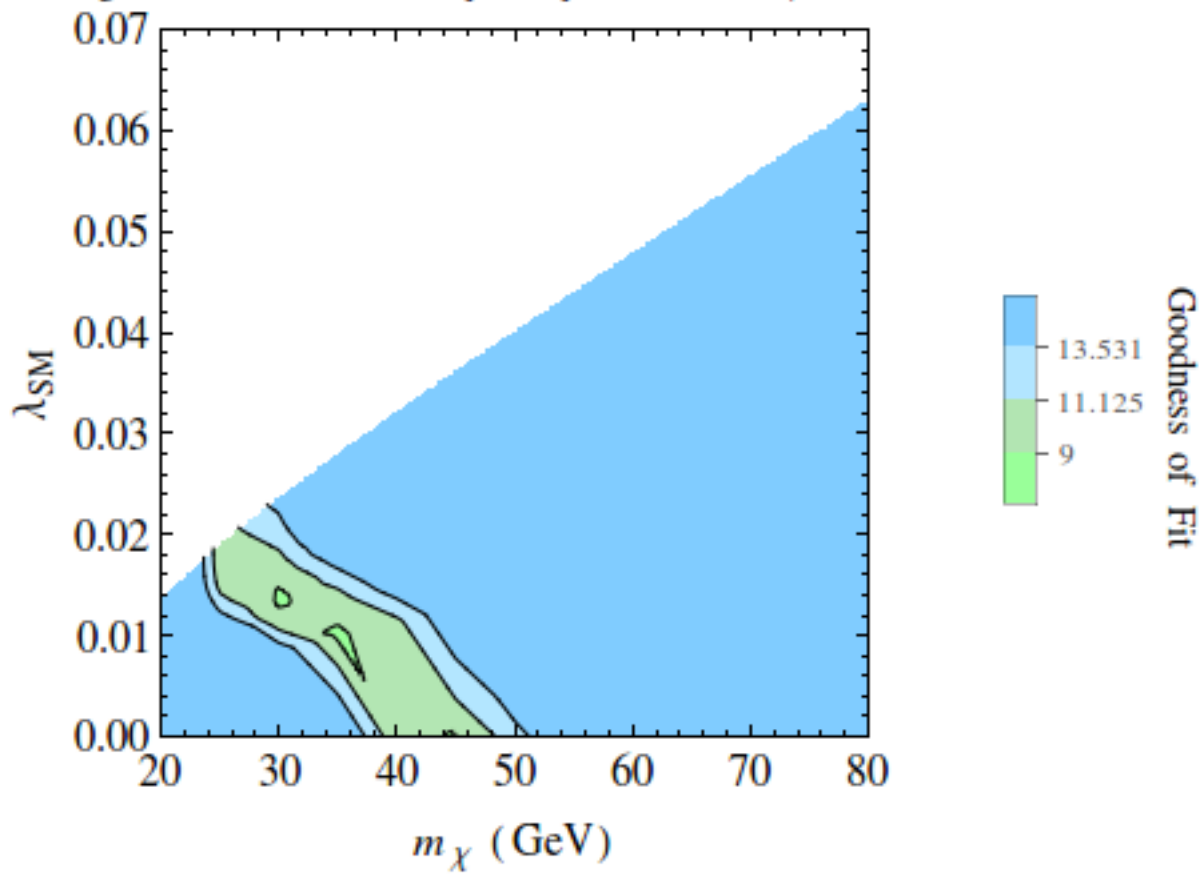
Significance of Fit for 6b Spectra – PseudoScalar

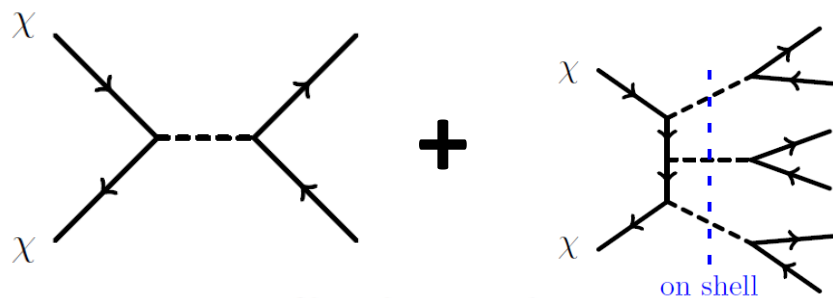


$$\lambda_{\text{DM}} \sim 1.1 - 1.4$$

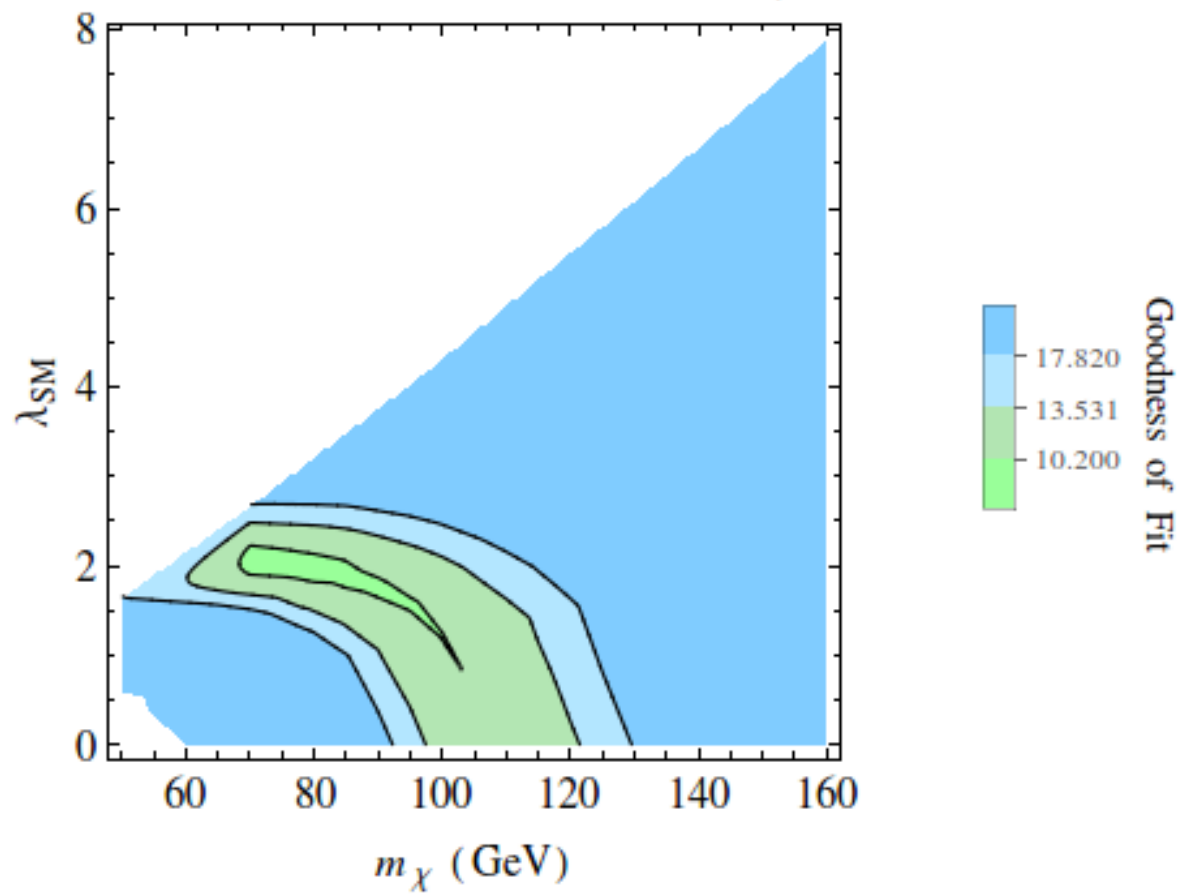


Significance of Fit for 4q + 2q - Vector, $m_V = 15$ GeV





Significance of Fit for 6b+2b – PseudoScalar, $m_\varphi = 20 \text{ GeV}$



Other On-shell mediator scenarios

- C. Boehm, M. J. Dolan and C. McCabe
arXiv:1404.4977 [hep-ph]
- P. Ko, W.-I. Park and Y. Tang
arXiv:1404.5257[hep-ph]
- A. Martin, J. Shelton and J. Unwin
arXiv:1405.0272 [hep-ph]

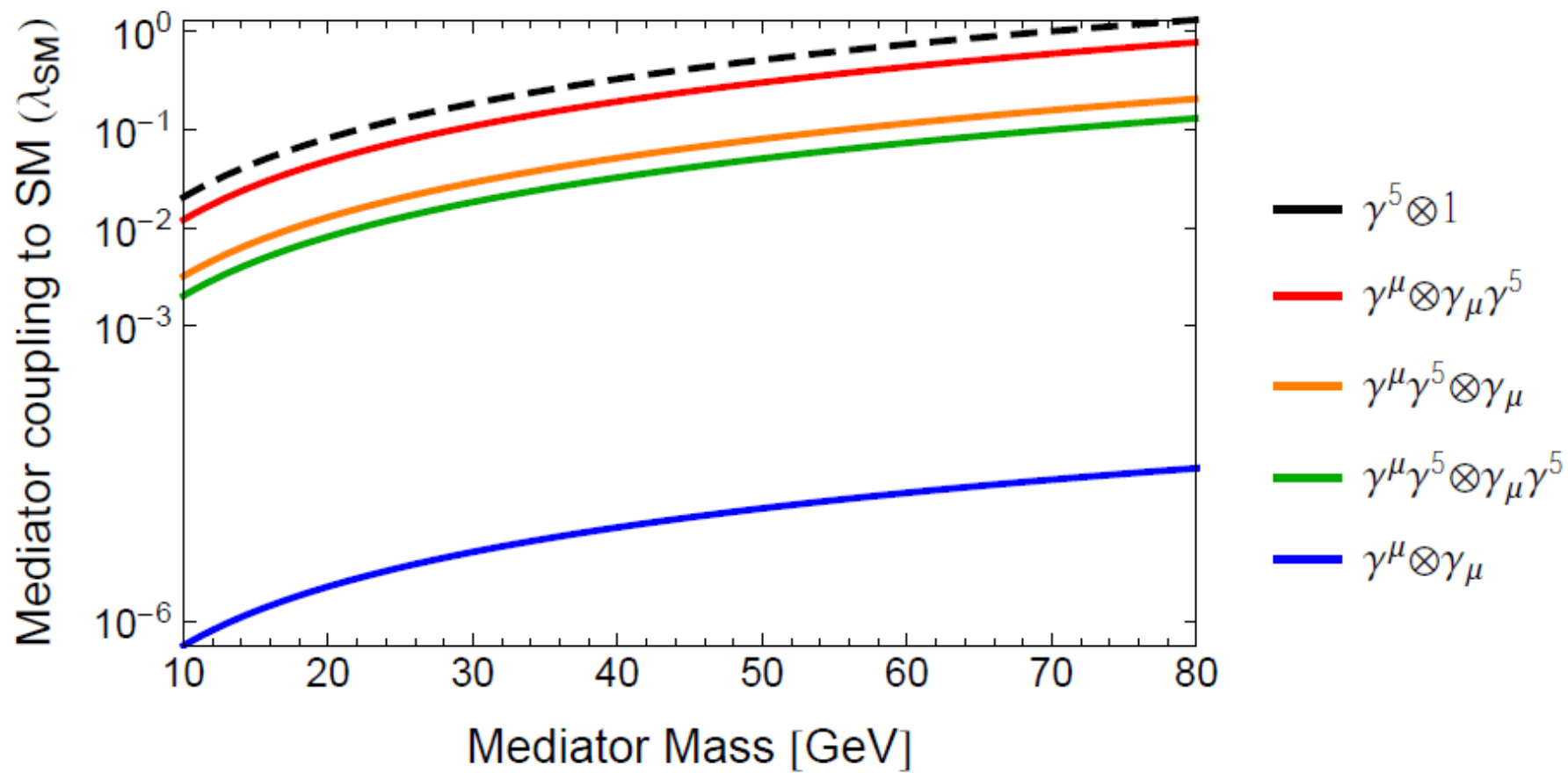
Bounds on hidden, on-shell mediators

Direct Detection

- Energy scale of direct detection experiments are on the order of 1-100 keV
- These energies allow the scattering to be parameterized as an effective theory where the relevant parameter is:

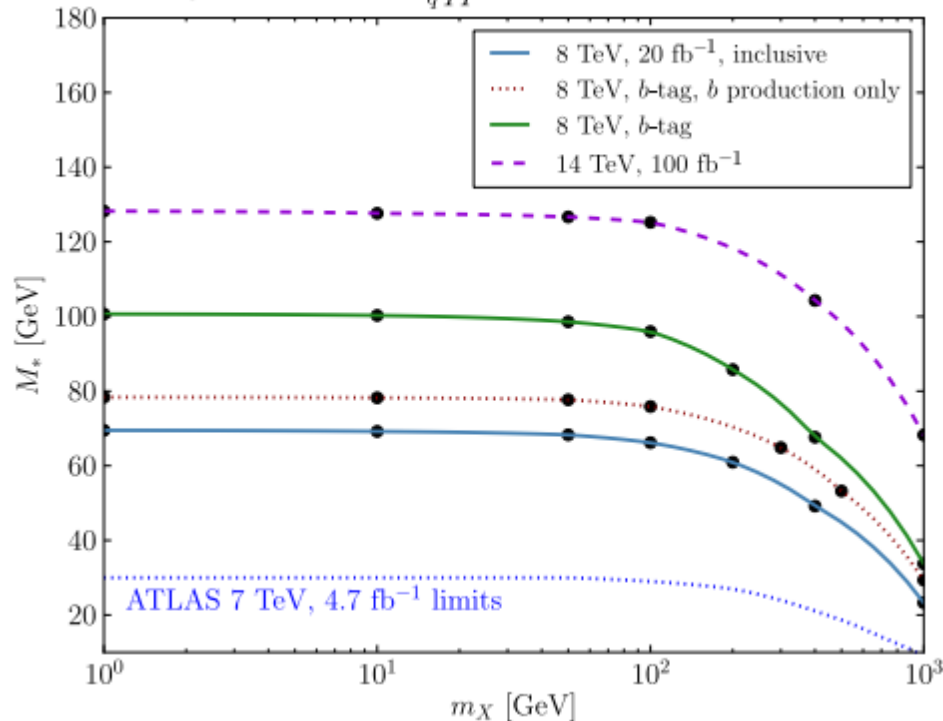
$$\frac{\lambda_{SM} \lambda_{DM}}{m_V^2}$$

$$\frac{\lambda_{SM} \lambda_{DM}}{m_\phi^2}$$



Colliders

Constraints on $m_q \bar{q} q \bar{X} X$ from mono- b searches



Collider Bounds

Projected bounds from
mono b searches

$$\lambda_{\text{SM}}^{\text{spin-0}} \lesssim 0.2$$

$$\lambda_{\text{SM}}^{\text{spin-1}} \lesssim 0.6$$

Opportunities for future work

- Relic Density too small

$$\frac{d\Phi(b, \ell)}{dE_\gamma} = \frac{\langle \sigma v \rangle_{b\bar{b}}}{2} \frac{1}{4\pi m_\chi^2} \frac{dN_\gamma}{dE_\gamma} \int_{\text{LOS}} dx \rho^2(r_{\text{gal}}(b, \ell, x))$$

One gets 2-3 times as many particles, but also dilutes the density by a factor of 2-3.

Conclusions

- On-shell mediators can be used to model the galactic center excess
 - WIMP mass 30-110 GeV
 - Excess can be explained with large amounts of freedom with regard to the standard model coupling
 - There are several experimental probes the standard model coupling

Backup Slides

